## PROBABILITY

# A KINGMAN CONVOLUTION APPROACH TO BESSEL PROCESSES* 

BY

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Dedicated to Professor Kazimierz Urbanik, my former teacher, from whom I learnt how a man and a mathematician could become to each other


#### Abstract

In this paper we study Bessel processes in terms of the Kingman convolution method. In particular, we propose a higher dimensional model of the Kingman convolution algebras. We show that every Bessel process started at 0 is induced by a Kingman convolution. Moreover, a new concept of increments of stochastic processes is introduced. It permits to regard Bessel processes as "stationary and independent increments processes".


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## 1. INTRODUCTION, NOTATION AND PRELIMILARIES

This study is inspired by a distinguished part of Bessel processes in financial mathematics for decades. Indeed, for each $n=1,2, \ldots$ let

$$
\mathbf{W}_{t}=\left(W_{t}^{(1)}, W_{t}^{(2)}, \ldots, W_{t}^{(n)}\right)
$$

be an $n$-dimensional Brownian motion $\left(B M^{(n)}\right)$ and $\rho_{t}=\left\|\mathbf{W}_{t}\right\|$ its radial part. Consider the following process:

$$
\begin{equation*}
\beta_{t}=\sum_{i=1}^{n} \int_{0}^{t} \frac{W_{s}^{(i)}}{\rho_{s}} d W_{s}^{(i)} \tag{1.1}
\end{equation*}
$$

[^0]which, since $\langle\beta, \beta\rangle_{t}=t$, stands for a linear Brownian motion, i.e. a $B M^{(1)}$. By virtue of Revuz and Yor [12], p. 439, we have
\[

$$
\begin{equation*}
\rho_{t}^{2}=\rho_{0}^{2}+2 \int_{0}^{t} \rho_{s} d \beta_{s}+n t \tag{1.2}
\end{equation*}
$$

\]

Replacing $n$ by any nonnegative number $\delta \geqslant 0$, we see that the equation (1.2) leads to the following interpolation class of stochastic differential equations (SDE):

$$
\begin{equation*}
Z_{t}=x+2 \int_{0}^{t} \sqrt{\left|Z_{s}\right|} d \beta_{s}+\delta t \tag{1.3}
\end{equation*}
$$

where $x \geqslant 0$.
Note that (1.3) is a special case of the Cox-Ingersoll-Ross family of diffusions [2] and has a unique solution which is strong, nonnegative and adapted with respect to the natural filtration $\left\{\mathcal{F}_{t}\right\}$ of $\left\{W_{t}\right\}$. Consequently, in the case when $\delta \geqslant 0$, $x \geqslant 0$, the absolute sign in (1.3) can be omitted and $\left\{Z_{t}\right\}$ can be modelled as short term interest rates (cf. Cox et al. [2]).

Definition 1.1 (cf. Revuz and Yor [12], XI). For every $\delta \geqslant 0, x \geqslant 0$, the unique strong solution of the equation (1.3) is called the square of $\delta$-dimensional Bessel process started at $x$ and is denoted by $B E S Q^{\delta}(x)$. Further, the square root of $B E S Q^{\delta}\left(x^{2}\right)$ is called the Bessel process ${ }^{1}$ of dimension $\delta$ started at $x$ and is denoted by $B E S^{\delta}(x)$.

In the sequel, we study the class of processes $B E S^{\delta}(x), \delta=2(s+1) \geqslant 1$, via the Kingman convolution method and also use $s$ as a fixed index of the Bessel process.

Let $\mathcal{P}$ denote the class of all probability measures (p.m.'s) on the positive halfline $\mathbb{R}^{+}$endowed with the weak convergence, and $*_{1, \delta}, \delta \geqslant 1$, denote the Kingman convolution which was introduced by Kingman [5] in connection with the addition of independent spherically symmetric random vectors (r.vec.'s) in a Euclidean space. Namely, for each continuous bounded function $f$ on $\mathbb{R}^{+}$we write
(1.4) $\int_{0}^{\infty} f(x) \mu *_{1, \delta} \nu(d x)$

$$
=\frac{\Gamma(s+1)}{\sqrt{\pi} \Gamma\left(s+\frac{1}{2}\right)} \int_{0}^{\infty} \int_{0}^{\infty} \int_{-1}^{1} f\left(\left(x^{2}+2 u x y+y^{2}\right)^{1 / 2}\right)\left(1-u^{2}\right)^{s-1 / 2} \mu(d x) \nu(d y) d u
$$

[^1]where $\mu, \nu \in P$ and $\delta=2(s+1) \geqslant 1$ (cf. Kingman [5] and Urbanik [15]). The algebra $\left(\mathcal{P}, *_{1, \delta}\right)$ is the most important example of Urbanik convolution algebras (cf. Urbanik [15]). In the language of the Urbanik convolution algebras, the characteristic measure, say $\sigma_{s}$, of the Kingman convolution has the Rayleigh density
\[

$$
\begin{equation*}
d \sigma_{s}(y)=\frac{2(s+1)^{s+1}}{\Gamma(s+1)} y^{2 s+1} \exp \left(-(s+1) y^{2}\right) d y \tag{1.5}
\end{equation*}
$$

\]

with the characteristic exponent $\varkappa=2$ and the kernel $\Lambda_{s}$,

$$
\begin{equation*}
\Lambda_{s}(x)=\Gamma(s+1) J_{s}(x) /(1 / 2 x)^{s} \tag{1.6}
\end{equation*}
$$

where $J_{s}(x)$ denotes the Bessel function,

$$
\begin{equation*}
J_{s}(x):=\sum_{k=0}^{\infty} \frac{(-1)^{k}(x / 2)^{\nu+2 k}}{k!\Gamma(\nu+k+1)} \tag{1.7}
\end{equation*}
$$

It is known (cf. Kingman [5], Theorem 1) that the kernel $\Lambda_{s}$ itself is an ordinary characteristic function (ch.f.) of a symmetric p.m., say $F_{s}$, defined on the interval $[-1,1]$. Thus, if $\theta_{s}$ denotes a random variable (r.v.) with distribution $F_{s}$, then for each $t \in \mathbb{R}^{+}$

$$
\begin{align*}
\Lambda_{s}(t) & =E \exp \left(i t \theta_{s}\right)  \tag{1.8}\\
& =\int_{-1}^{1} \exp (i t x) d F_{s}(x)
\end{align*}
$$

Suppose that $X$ is a nonnegative r.v. with distribution $\mu \in \mathcal{P}$ and $X$ is independent of $\theta_{s}$. The radial characteristic function (rad.ch.f.) of $\mu$, denoted by $\hat{\mu}(t)$, is defined by

$$
\begin{align*}
\hat{\mu}(t) & =E \exp \left(i t X \theta_{s}\right)  \tag{1.9}\\
& =\int_{0}^{\infty} \Lambda_{s}(t x) \mu(d x)
\end{align*}
$$

for every $t \in \mathbb{R}^{+}$. In particular, the rad.ch.f. of $\sigma_{s}$ is

$$
\begin{equation*}
\hat{\sigma}_{s}(t)=\exp \left(-\frac{t^{2}}{2}\right), \quad t \in \mathbb{R}^{+} \tag{1.10}
\end{equation*}
$$

It should be noted, since the rad.ch.f. is defined uniquely up to the mapping $x \rightarrow a x, a>0, x \in \mathbb{R}^{+}$, that the representation (1.10) may differ from the one given in Urbanik [15] and Kingman [5] only by a scale parameter.

## 2. CARTESIAN PRODUCT OF KINGMAN CONVOLUTIONS

Denote by $\mathbb{R}^{+k}, k=1,2, \ldots$, the $k$-dimensional nonnegative cone of $\mathbb{R}^{k}$ and $\mathcal{P}\left(\mathbb{R}^{+k}\right)$ the class of all p.m.'s on $\mathbb{R}^{+k}$ equipped with the weak convergence. In the sequel, we will denote the multidimensional vectors and distributions and r.vec.'s by bold letters. For each point $z$ of any set $Z$ let $\delta_{z}$ denote the Dirac measure (the unit mass) at the point $z$. In particular, if $\mathbf{x}=\left(x_{1}, x_{2}, \ldots, x_{k}\right) \in \mathbb{R}^{k+}$, then

$$
\begin{equation*}
\delta_{\mathbf{x}}=\delta_{x_{1}} \times \delta_{x_{2}} \times \ldots \times \delta_{x_{k}}, \tag{2.1}
\end{equation*}
$$

where the sign $\times$ denotes the Cartesian product of measures. We put, for $\mathbf{x}=$ $\left(x_{1}, x_{2}, \ldots, x_{k}\right)$ and $\mathbf{y}=\left(y_{1}, y_{2}, \ldots, y_{k}\right) \in \mathbb{R}^{+k}$,

$$
\begin{equation*}
\delta_{\mathbf{x}} \bigcirc_{k} \delta_{\mathbf{y}}=\left\{\delta_{x_{1}} \circ \delta_{y_{1}}\right\} \times\left\{\delta_{x_{2}} \circ \delta_{y_{2}}\right\} \times \ldots \times\left\{\delta_{x_{k}} \circ \delta_{y_{k}}\right\} ; \tag{2.2}
\end{equation*}
$$

for the sake of simplicity, here and somewhere below we denote the Kingman convolution operation $*_{1, \delta}$ simply by o . Since convex combinations of p.m.'s of the form (2.1) are dense in $\mathcal{P}\left(\mathbb{R}^{+k}\right)$, the relation (2.2) can be extended to arbitrary p.m.'s $\mathbf{F}, \mathbf{G} \in \mathcal{P}\left(\mathbb{R}^{+k}\right)$. Namely, we put

$$
\begin{equation*}
\mathbf{F} \bigcirc_{k} \mathbf{G}=\iint_{\mathbb{R}^{+k}} \delta_{\mathbf{x}} \bigcirc_{k} \delta_{\mathbf{y}} \mathbf{F}(d \mathbf{x}) \mathbf{G}(d \mathbf{y}) . \tag{2.3}
\end{equation*}
$$

In the sequel, the binary operation $\bigcirc_{k}$ will be called the $k$-times Cartesian product of Kingman convolutions. It is easy to show that the binary operation $\mathrm{O}_{k}$ is continuous in the weak topology, which together with (1.4) and (2.3) implies the following theorem.

Theorem 2.1. The pair $\left(\mathcal{P}\left(\mathbb{R}^{+k}\right), O_{k}\right)$ is a commutative topological semigroup with $\delta_{0}$ as the unit element. Moreover, the operation $\mathrm{O}_{k}$ is distributive with respect to convex combinations of p.m.'s in $\mathcal{P}\left(\mathbb{R}^{+k}\right)$.

In the sequel, the pair $\left(\mathcal{P}\left(\mathbb{R}^{+k}\right), \bigcirc_{k}\right)$ will be called a $k$-dimensional Kingman convolution algebra ${ }^{2}$. For every $\mathbf{F} \in \mathcal{P}\left(\mathbb{R}^{+k}\right)$ the $k$-dimensional rad.ch.f. $\hat{\mathbf{F}}(\mathbf{t}), \mathbf{t}=\left(t_{1}, t_{2}, \ldots, t_{k}\right) \in \mathbb{R}^{k+}$, is defined by

$$
\begin{equation*}
\hat{\mathbf{F}}(\mathbf{t})=\int_{\mathbb{R}^{+k}} \prod_{j=1}^{k} \Lambda_{s}\left(t_{j} x_{j}\right) \mathbf{F}(\mathbf{d x}), \tag{2.4}
\end{equation*}
$$

where $\mathbf{x}=\left(x_{1}, x_{2}, \ldots, x_{k}\right) \in \mathbb{R}^{+k}$.
The $k$-dimensional Rayleigh distribution, say $\boldsymbol{\Sigma}_{s}$, is defined by

$$
\begin{equation*}
\boldsymbol{\Sigma}_{s}=\sigma_{s} \times \sigma_{s} \times \ldots \times \sigma_{s} \quad(k \text { times }) . \tag{2.5}
\end{equation*}
$$

[^2]Furthermore, for any nonnegative numbers $\lambda_{r}, r=1,2, \ldots$, the distribution

$$
\begin{equation*}
\mathbf{F}=\left\{T_{\lambda_{1}} \sigma_{s}\right\} \times\left\{T_{\lambda_{2}} \sigma_{s}\right\} \times \ldots \times\left\{T_{\lambda_{k}} \sigma_{s}\right\} \tag{2.6}
\end{equation*}
$$

stands for a $k$-dimensional Rayleighian distribution. Here and in the sequel, if $X$ is an r.v. or an r.vec. with distribution $\mu$ and $\lambda$ is a real number, then we denote by $T_{\lambda} \mu$ the distribution of $\lambda X$.

By virtue of formulas (1.10) and (2.4)-(2.6) we have the following
THEOREM 2.2. Suppose distributions $\boldsymbol{\Sigma}$ and $\mathbf{F}$ are of the form (2.5) and (2.6). Then, for any $\mathbf{t} \in \mathbb{R}^{+k}$,

$$
\begin{equation*}
-\log \hat{\boldsymbol{\Sigma}}_{s}(\mathbf{t})=\frac{1}{2} \sum_{j=1}^{k} t_{j}^{2} \tag{2.7}
\end{equation*}
$$

and

$$
\begin{equation*}
-\log \hat{\mathbf{F}}(\mathbf{t})=\frac{1}{2} \sum_{j=1}^{k} \lambda_{j}^{2} t_{j}^{2} \tag{2.8}
\end{equation*}
$$

Let $\theta, \theta_{1}, \theta_{2}, \ldots, \theta_{k}$ be independent identically distributed (i.i.d.) r.v.'s with common distribution $F_{s}$. We set

$$
\begin{equation*}
\boldsymbol{\Theta}_{s}=\left(\theta_{1}, \theta_{2}, \ldots, \theta_{k}\right) \tag{2.9}
\end{equation*}
$$

Assume that $\mathbf{X}=\left(X_{1}, X_{2}, \ldots, X_{k}\right)$ is a $k$-dimensional r.vec. with distribution $\mathbf{F}$ and $\mathbf{X}$ is independent of $\boldsymbol{\Theta}$. We put

$$
\begin{equation*}
[\boldsymbol{\Theta}, \mathbf{X}]=\left(\theta_{1} X_{1}, \theta_{2} X_{2}, \ldots, \theta_{k} X_{k}\right) \tag{2.10}
\end{equation*}
$$

Then the following formula is the multidimensional generalization of (1.9) and is equivalent to (2.4):

$$
\begin{equation*}
\hat{\mathbf{F}}(\mathbf{t})=E \exp (i\langle\mathbf{t},[\boldsymbol{\Theta}, \mathbf{X}]\rangle) \tag{2.11}
\end{equation*}
$$

where $\mathbf{X}$ and $\boldsymbol{\Theta}$ are assumed to be independent, $\mathbf{t}=\left(t_{1}, t_{2}, \ldots, t_{k}\right) \in \mathbb{R}^{+k}$, and the symbol $\langle$,$\rangle denotes the inner product in \mathbb{R}^{k}$. In fact, we have
(2.12) $E \exp \left(i\left\langle\left(\theta_{1} t_{1}, \theta_{2} t_{2}, \ldots, \theta_{k} t_{k}\right), \mathbf{X}\right\rangle\right)=\int_{\mathbb{R}^{+k}} E \exp \left(i \sum_{j=1}^{k} t_{j} x_{j} \theta_{j}\right) \mathbf{F}(d \mathbf{x})$

$$
=\int_{\mathbb{R}^{+k}} \prod_{j=1}^{k} \Lambda_{s}\left(t_{j} x_{j}\right) \mathbf{F}(d \mathbf{x})=\hat{\mathbf{F}}(\mathbf{t}) .
$$

As a consequence of the representation (2.11) we have
COROLLARY 2.1. For each $\mathbf{F} \in \mathcal{P}\left(\mathbb{R}^{k+}\right)$ the rad.ch.f. $\hat{\mathbf{F}}(\mathbf{t})$ is also an ordinary $k$-dimensional ch.f., and hence it is uniformly continuous.

The following lemma will be used in the representation of $k$-dimensional infinitely divisible (ID) p.m.'s.

Lemma 2.1. (i) For every $t \geqslant 0$

$$
\begin{equation*}
\lim _{x \rightarrow 0} \frac{1-\Lambda_{s}(t x)}{x^{2}}=\lim _{x \rightarrow 0} \frac{1-E e^{i t x \theta}}{x^{2}}=\frac{t^{2}}{2} \tag{2.13}
\end{equation*}
$$

(ii) For any vectors $\mathbf{x}=\left(x_{0}, x_{1}, \ldots, x_{k}\right)$ and $\mathbf{t}=\left(t_{0}, t_{1}, \ldots, t_{k}\right) \in \mathbb{R}^{k+1}$, $k=1,2, \ldots$,

$$
\begin{equation*}
\lim _{\rho \rightarrow 0} \frac{1-\prod_{r=0}^{k} \Lambda_{s}\left(t_{r} x_{r}\right)}{\rho^{2}}=\sum_{r=0}^{k} \lambda_{r}(\mathbf{x}) t_{r}^{2} \tag{2.14}
\end{equation*}
$$

with $\rho=\|\mathbf{x}\|$ and $\lambda_{r}(\mathbf{x}), r=0,1, \ldots, k$, given by

$$
\lambda_{r}(\mathbf{x})= \begin{cases}\frac{1}{2} \cos ^{2} \phi & r=0  \tag{2.15}\\ \frac{1}{2}\left(\sin \phi \sin \phi_{1} \ldots \sin \phi_{r-1} \cos \phi_{r}\right)^{2}, & 1 \leqslant r \leqslant k-2 \\ \frac{1}{2}\left(\sin \phi \sin \phi_{1} \ldots \sin \phi_{k-2} \cos \psi\right)^{2}, & r=k-1 \\ \frac{1}{2}\left(\sin \phi \sin \phi_{1} \ldots \sin \phi_{k-2} \sin \psi\right)^{2}, & r=k\end{cases}
$$

where $0 \leqslant \psi, \phi, \phi_{r} \leqslant \pi / 2, r=1,2, \ldots, k-2$, are angles of x appearing in its polar form.

Proof. (i) The equation (1.8) together with the l'Hôpital rule implies that

$$
\lim _{x \rightarrow 0} \frac{1-\Lambda_{s}(t x)}{x^{2}}=\lim _{x \rightarrow 0} \frac{1-E e^{i t \theta}}{x^{2}}=\frac{t^{2}}{2}
$$

which proves (2.13).
(ii) In order to prove (2.14) assume that the points $\mathbf{x}=\left(x_{0}, x_{1}, \ldots, x_{k}\right) \in$ $\mathbb{R}^{k+1}$ are of the polar form

$$
x_{r}= \begin{cases}\rho \cos \phi, & r=0  \tag{2.16}\\ \rho \sin \phi \sin \phi_{1} \ldots \sin \phi_{r-1} \cos \phi_{r}, & 1 \leqslant r \leqslant k-2 \\ \rho \sin \phi \sin \phi_{1} \ldots \sin \phi_{k-2} \cos \psi, & r=k-1 \\ \rho \sin \phi \sin \phi_{1} \ldots \sin \phi_{k-2} \sin \psi, & r=k\end{cases}
$$

where $0 \leqslant \psi, \phi, \phi_{r} \leqslant \pi / 2, r=1,2, \ldots, k-2$. Putting
(2.17) $A(\boldsymbol{\Theta}, \mathbf{t}, \boldsymbol{\Phi})=t_{0} \theta_{0} \cos \phi+\sum_{r=1}^{k-2} t_{r} \theta_{r} \sin \phi \sin \phi_{1} \ldots \sin \phi_{r-1} \cos \phi_{r}$ $+t_{k-1} \theta_{k-1} \sin \phi \sin \phi_{1} \ldots \sin \phi_{k-2} \cos \psi+t_{k} \theta_{k} \sin \phi \sin \phi_{1} \ldots \sin \phi_{k-2} \sin \psi$
and

$$
\begin{equation*}
V(\boldsymbol{\Theta}, \mathbf{t}, \boldsymbol{\Phi})=\sum_{r=0}^{k} t_{r} x_{r} \theta_{r} \tag{2.18}
\end{equation*}
$$

where the $\theta_{r}, r=0,1,2, \ldots$, are symmetric i.i.d. r.v.'s with distribution $\sigma_{s}, \boldsymbol{\Phi}=$ $\left(\psi, \phi, \phi_{1}, \ldots, \phi_{k}\right)$ and $\boldsymbol{\Theta}:=\left(\theta_{0}, \theta_{1}, \ldots, \theta_{k}\right)$. By virtue of (2.12) and (2.16) and applying l'Hôpital rule, we have

$$
\text { (2.19) } \begin{aligned}
\lim _{\rho \rightarrow 0} \frac{1-\prod_{r=0}^{k} \Lambda_{s}\left(t_{r} x_{r}\right)}{\rho^{2}} & =\lim _{\rho \rightarrow 0} \frac{1-E\left(\exp \left(i \sum_{r=0}^{k} t_{r} x_{r} \theta_{r}\right)\right)}{\rho^{2}} \\
& =\left.\frac{\left(d^{2} / d \rho^{2}\right)(1-E \exp (i \rho A(\boldsymbol{\Theta}, \mathbf{t}, \boldsymbol{\Phi})))}{\left(d^{2} / d \rho^{2}\right) \rho^{2}}\right|_{\rho=0} \\
& =\left.\frac{1}{2} E V^{2}(\boldsymbol{\Theta}, \mathbf{t}, \mathbf{\Phi}) \exp (i \rho V(\boldsymbol{\Theta}, \mathbf{t}, \mathbf{\Phi}))\right|_{\rho=0}
\end{aligned}
$$

Since $\sigma_{s}$ has expectation zero and variance one, it follows that

$$
\begin{equation*}
E V^{2}(\theta, \mathbf{t}, \phi)=\sum_{j=1}^{k} t_{j}^{2} x_{j}^{2} \tag{2.20}
\end{equation*}
$$

which together with (2.19) implies (2.14).
Proceeding successively, we have the following theorem:
THEOREM 2.3. Every p.m. $\mathbf{F} \in \mathcal{P}\left(\mathbb{R}^{+k}\right)$ is uniquely determined by its $k$ dimensional rad.ch.f. $\hat{\mathbf{F}}$ and the following formula holds:

$$
\begin{equation*}
\widehat{\mathbf{F}}_{1}{\widehat{O_{k}} \mathbf{F}}_{2}(\mathbf{t})=\widehat{\mathbf{F}_{1}}(\mathbf{t}) \widehat{\mathbf{F}_{2}}(\mathbf{t}) \tag{2.21}
\end{equation*}
$$

where $\mathbf{F}_{1}, \mathbf{F}_{2} \in \mathcal{P}\left(\mathbb{R}^{+k}\right)$ and $\mathbf{t} \in \mathbb{R}^{+k}$.
Proof. The formula (2.21) follows from (1.4) and (2.3). Next, using the formulas (2.3) and (2.4) and integrating the function $\hat{\mathbf{F}}\left(t_{1} u_{1}, \ldots, t_{k} u_{k}\right) k$ times with respect to $\sigma_{s}$, we get

$$
\begin{align*}
\int_{\mathbb{R}^{+k}} \hat{\mathbf{F}}\left(t_{1} u_{1}, \ldots,\right. & \left.t_{k} u_{k}\right) \sigma_{s}\left(d u_{1}\right) \ldots \sigma_{s}\left(d u_{k}\right)  \tag{2.22}\\
& =\int_{\mathbb{R}^{+}} \ldots \int_{\mathbb{R}^{+}} \prod_{j=1}^{k} \Lambda_{s}\left(t_{j} x_{j} u_{j}\right) \mathbf{F}(\mathbf{d} \mathbf{x}) \sigma_{s}\left(d u_{1}\right) \ldots \sigma_{s}\left(d u_{k}\right) \\
& =\int_{\mathbb{R}^{+k}} \prod_{j=1}^{k} \exp \left(-t_{j}^{2} x_{j}^{2}\right) \mathbf{F}(\mathbf{d} \mathbf{x})
\end{align*}
$$

which, by the change of variables $y_{j}=x_{j}^{2}, j=1, \ldots, k$, and by the uniqueness of the $k$-dimensional Laplace transform, implies that $\mathbf{F}$ is uniquely determined by the left-hand side of (2.22).

As a consequence of the formula (2.22) we have the following corollary which is an analogue of the continuity theorem for multidimensional Laplace transforms.

THEOREM 2.4. Suppose that $\left\{\mathbf{F}_{n}\right\}$ is a sequence of distributions on $\mathbb{R}^{k+}$ and $\left\{\phi_{n}\right\}$ is a sequence of the corresponding rad.ch.f.'s. Then $\mathbf{F}_{n}$ converges weakly to a distribution $\mathbf{F}$ if and only if $\left\{\phi_{n}\right\}$ converges uniformly on every compact subset of $\mathbb{R}^{k+}$ to a rad.ch.f. $\phi$.

For any $\mathbf{x} \in \mathbb{R}^{+k}$ the generalized translation operators (g.t.o.'s) $\mathbf{T}^{\mathbf{x}}$ acting on the Banach space $\mathbb{C}_{b}\left(\mathbb{R}^{+k}\right)$ of real bounded continuous functions $f$ on $\mathbb{R}^{+k}$ are defined, for each $\mathbf{y} \in \mathbb{R}^{+k}$, by

$$
\begin{equation*}
\mathbf{T}^{\mathbf{x}} f(\mathbf{y})=\int_{\mathbb{R}^{+k}} f(\mathbf{u})\left\{\delta_{\mathbf{x}} \bigcirc_{k} \delta_{\mathbf{y}}\right\}(d \mathbf{u}) \tag{2.23}
\end{equation*}
$$

In terms of these g.t.o.'s the $k$-dimensional rad.ch.f. of p.m.'s on $\mathbb{R}^{+k}$ can be characterized as follows (see Vólkovich [17] for the proof):

THEOREM 2.5. A real bounded continuous function $f$ on $\mathbb{R}^{+k}$ is a ( $k$-dimensional) rad.ch.f. of a p.m. if and only if $f(\mathbf{0})=1$ andf is $\left\{\mathbf{T}^{\mathbf{x}}\right\}$-nonnegative definite in the sense that for any $\mathbf{x}_{1}, \mathbf{x}_{2}, \ldots, \mathbf{x}_{k} \in \mathbb{R}^{k}$ and $\lambda_{1}, \lambda_{2}, \ldots, \lambda_{k} \in \mathbb{C}$

$$
\begin{equation*}
\sum_{i, j=1}^{k} \lambda_{i} \bar{\lambda}_{j} \mathbf{T}^{\mathbf{x}_{i}} f\left(\mathbf{x}_{j}\right) \geqslant 0 \tag{2.24}
\end{equation*}
$$

The $k$-dimensional ID elements with respect to $\bigcirc_{k}$ can be defined as follows:
Definition 2.1. A p.m. $\mu \in \mathcal{P}\left(\mathbb{R}^{+k}\right)$ is called infinitely divisible (ID) if for every natural $m$ there exists a p.m. $\mu_{m}$ such that

$$
\mu=\mu_{m} \bigcirc_{k} \ldots \bigcirc_{k} \mu_{m} \quad(m \text { times })
$$

The simplest but most important example of $k$-dimensional ID distributions are the $k$-dimensional Rayleigh distributions. More generally, if $\mathbf{F}$ is a $k$-dimensional Rayleighian distribution, then it is also ID. Let us denote by $I D\left(\bigcirc_{k}\right)$ the class of all i.d.p.m.'s in $\left(\mathcal{P}\left(\mathbb{R}^{+k}\right), \bigcirc_{k}\right)$. The following theorem, being a generalization of Theorem 7 in Kingman [5], stands for an analogue of the Lévy-Khintchine representation for rad.ch.f.'s of i.d.p.m.'s in the $k$-dimensional Kingman convolution.

THEOREM 2.6. A p.m. $\mu \in I D\left(\bigcirc_{k}\right)$ if and only if there exists a $\sigma$-finite measure $M$ (a Lévy measure) on $\mathbb{R}^{+k}$ with the property that $M(\{\mathbf{0}\})=0, M$ is finite
outside every neighborhood of $\mathbf{0}$ and

$$
\begin{equation*}
\int_{\mathbb{R}^{+k}} \frac{\|\mathbf{x}\|^{2}}{1+\|\mathbf{x}\|^{2}} M(d \mathbf{x})<\infty \tag{2.25}
\end{equation*}
$$

and, for each $\mathbf{t}=\left(t_{1}, \ldots, t_{k}\right) \in \mathbb{R}^{k+}$,

$$
\begin{equation*}
-\log \hat{\mu}(\mathbf{t})=\int_{\mathbb{R}^{+k}}\left\{1-\prod_{j=1}^{k} \Lambda_{s}\left(t_{j} x_{j}\right)\right\} \frac{1+\|\mathbf{x}\|^{2}}{\|\mathbf{x}\|^{2}} M(d \mathbf{x}) \tag{2.26}
\end{equation*}
$$

where, at the origin $\mathbf{0}$, the integrand on the right-hand side of (2.26) is assumed to be of the form

$$
\begin{equation*}
\sum_{j=1}^{k} \lambda_{j}(\mathbf{x}) t_{j}^{2}=\lim _{\|\mathbf{x}\| \rightarrow 0}\left\{1-\prod_{j=1}^{k} \Lambda_{s}\left(t_{j} x_{j}\right)\right\} \frac{1+\|\mathbf{x}\|^{2}}{\|\mathbf{x}\|^{2}} \tag{2.27}
\end{equation*}
$$

for nonnegative $\lambda_{j}(\mathbf{x}), j=1,2, \ldots, k$, and $\mathbf{x} \in \mathbb{R}^{k+}$, given by equations (2.15). In particular, if $M$ tends to the measure $\mathbf{0}$, then $\mu$ becomes a Rayleighian distribution with the rad.ch.f.

$$
\begin{equation*}
-\log \hat{\mu}(\mathbf{t})=\frac{1}{2} \sum_{j=1}^{k} \lambda_{j} t_{j}^{2}, \quad \mathbf{t} \in \mathbb{R}^{k+} \tag{2.28}
\end{equation*}
$$

for some nonnegative $\lambda_{j}, j=1, \ldots, k$.
Moreover, the representation (2.26) is unique.
Proof. The proof is carried out in several steps.
(i) If $\phi$ is a $k$-dimensional ID rad.ch.f., then it does not vanish on $\mathbb{R}^{k+}$.

Indeed, denote by $\Phi_{k}$ the totality of $k$-dimensional ID rad.ch.f.'s (of the fixed index $s$ ). Then, we have

$$
\begin{equation*}
\Phi_{k}=\bigcap_{n=1}^{\infty}\left\{\phi: \phi^{1 / n} \in \Phi_{n}\right\} \tag{2.29}
\end{equation*}
$$

which, together with (2.12) and (2.21), implies that every $k$-dimensional ID rad.ch.f. is a symmetric ordinary ID ch.f. and, consequently, it does not vanish on $\mathbb{R}^{k+}$.
(ii) Any $\nu \in I D\left(\bigcirc_{k}\right)$ with rad.ch.f. $\hat{\nu}=\psi \in \Phi_{k}$ can be expressed in the form (2.26).

Accordingly, for every $n$ there exists $\psi_{n} \in \Phi_{k}$ such that $\psi=\psi_{n}^{n}$. By virtue of (i), $\psi(\mathbf{t})>0$ for each $\mathbf{t}$. Therefore,

$$
\begin{equation*}
\log \psi(\mathbf{t})=\lim _{n \rightarrow \infty} n\left\{\psi_{n}(\mathbf{t})-1\right\} \tag{2.30}
\end{equation*}
$$

Let $H_{n}$ be a p.m. such that

$$
\begin{equation*}
\psi_{n}(\mathbf{t})=\int_{\mathbb{R}^{k+}} \prod_{j=1}^{k} \Lambda_{s}\left(t_{j} x_{j}\right) \mathbf{H}_{n}(d \mathbf{x}), \quad \mathbf{t} \in \mathbb{R}^{k+} \tag{2.31}
\end{equation*}
$$

Putting

$$
\begin{equation*}
\mathbf{G}_{n}(A)=n \int_{A} \frac{\|\mathbf{x}\|^{2}}{1+\|\mathbf{x}\|^{2}} \mathbf{H}_{n}(d \mathbf{x}) \tag{2.32}
\end{equation*}
$$

and taking into account the equations (2.30) and (2.31) we get

$$
\begin{equation*}
-\log \psi(\mathbf{t})=\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{k+}}\left\{1-\prod_{j=1}^{k} \Lambda_{s}\left(t_{j} x_{j}\right)\right\} \frac{1+\|\mathbf{x}\|^{2}}{\|\mathbf{x}\|^{2}} \mathbf{G}_{n}(d \mathbf{x}) \tag{2.33}
\end{equation*}
$$

which can be rewritten as

$$
\begin{equation*}
-\log \psi(\mathbf{t})=\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{k+}}\left\{1-\prod_{j=1}^{k} \Lambda_{s}\left(t_{j} x_{j}\right)\right\} \mathbf{K}_{n}(d \mathbf{x}) \tag{2.34}
\end{equation*}
$$

where $\mathbf{K}_{n}$ are finite measures vanishing at $\mathbf{0}$ defined by

$$
\mathbf{K}_{n}(d \mathbf{x}):=\frac{1+\|\mathbf{x}\|^{2}}{\|\mathbf{x}\|^{2}} \mathbf{G}_{n}(d \mathbf{x}) \quad(n=1,2, \ldots)
$$

Replacing $\mathbf{t}$ in the equation (2.34) by $[\mathbf{t}, \mathbf{u}], \mathbf{t}, \mathbf{u} \in \mathbb{R}^{k+}$, and integrating with respect to $\sigma_{s} \times \ldots \times \sigma_{s}(d \mathbf{u})$, we obtain

$$
\begin{aligned}
& -\int_{\mathbb{R}^{k+}} \log \psi([\mathbf{t}, \mathbf{u}]) \sigma_{s} \times \ldots \times \sigma_{s}(d \mathbf{u}) \\
& \quad=\int_{\mathbb{R}^{k+}} \lim _{n \rightarrow \infty} \int_{\mathbb{R}^{k+}}\left\{1-\prod_{j=1}^{k} \Lambda_{s}\left(t_{j} u_{j} x_{j}\right)\right\} \mathbf{K}_{n}(d \mathbf{x}) \sigma_{s} \times \ldots \times \sigma_{s}(d \mathbf{u}) \\
& \quad=\lim _{n \rightarrow \infty} \int_{\mathbb{R}^{k+}}\left\{1-\prod_{j=1}^{k} \exp \left(-t_{j}^{2} x_{j}^{2}\right)\right\} \mathbf{K}_{n}(d \mathbf{x})
\end{aligned}
$$

which, by changing variables $x_{j}^{2} \rightarrow u_{j}, j=1,2, \ldots, k$, and applying the continuity theorem for the classical infinitely divisible Laplace transforms on $\mathbb{R}^{k+}$, implies that there exists a finite measure $\mathbf{K}$ vanishing at $\mathbf{0}$ and a subsequence $\left\{\mathbf{K}_{m_{r}}\right\}$ which converges to $\mathbf{K}$ in the sense that for any bounded continuous function $f$ from $\mathbb{R}^{k+}$ to $\mathbb{R}$ vanishing on a neighborhood of $\mathbf{0}$ and

$$
\lim _{r \rightarrow \infty} \int_{\mathbb{R}^{k+}} f(\mathbf{x}) \mathbf{K}_{m_{r}}(d \mathbf{x})=\int_{\mathbb{R}^{k+}} f(\mathbf{x}) \mathbf{K}(d \mathbf{x})
$$

This together with (2.33) and (2.14) implies that every $\psi$ is of the form (2.26) for a Lévy measure $\mathbf{M}$.
(iii) Now, if $\mathbf{M}$ tends to the zero measure, it follows that, at the origin $\mathbf{0}$, the integrand on the right-hand side of (2.26) is determined by (2.1), which is a consequence of Lemma 2.1.
(iv) Conversely, the uniqueness of the formula (2.26) can be proved in the same way as in the classical case (cf. Sato [13], Theorems 8.1 and 8.7).

## 3. CONVOLUTION STRUCTURE OF BESSEL PROCESSES

Given a p.m. $\mu \in \mathcal{P}$ and $n=1,2, \ldots$ we put, for any $x \in \mathbb{R}^{+}, B \in \mathcal{B}\left(\mathbb{R}^{+}\right)$, where $\mathcal{B}\left(\mathbb{R}^{+}\right)$denotes the Borel $\sigma$-field of $\mathbb{R}^{+}$,

$$
\begin{equation*}
P_{n}(x, E)=\delta_{x} \circ \mu^{\circ n}(E) \tag{3.1}
\end{equation*}
$$

where the power is taken in the convolution o sense. Using the rad.ch.f. one can show that $\left\{P_{n}(x, E)\right\}$ satisfies the Chapman-Kolmogorov equation, and therefore there exists a nonnegative homogeneous Markov sequence, say $\left\{S_{n}^{x}\right\}, n=$ $0,1,2, \ldots$, with transition probability $\left\{P_{n}(x, E)\right\}$.

In what follows we will discuss the case of Bessel processes which stand for a continuous counterpart of the above symmetric random walks. Namely, suppose that $\mu$ is ID with respect to the Kingman convolution $\circ$. We put

$$
\begin{equation*}
q(t, x, E):=\mu^{\circ t} \circ \delta_{x}(E) \tag{3.2}
\end{equation*}
$$

and take into account the fact that the family $\{q(t, x, \cdot)\}$ of distributions satisfies the Chapman-Kolmogorov equation, and therefore it stands for a transition probability of a homogeneous strong Markov Feller process, say $\left\{X_{t}^{x}\right\}, t, x \in \mathbb{R}^{+}$, and, moreover, $\left\{X_{t}^{x}\right\}$ is stochastically continuous and has a cadlag version (cf. Nguyen [7], Theorem 2.6).

Definition 3.1. A stochastic process $\left\{X_{t}^{x}\right\}$ is called a Lévy-type (or o-Lévy) process if
(i) $X_{0}^{x}=x$ with probability 1 ;
(ii) $\left\{X_{t}^{x}\right\}$ is a strong Markov Feller process with transition probability of the form (3.2);
(iii) $\left\{X_{t}^{x}\right\}$ is a stochastically continuous process having cadlag realizations with probability 1.

It is evident that all Lévy processes are $*$-Lévy ones. The simplest example of Lévy-type but non-Lévy processes is the absolute value of the linear Brownian motion. Similarly, the following theorem shows that Bessel processes started from 0 stand for Lévy-type processes induced by the Kingman convolution.

THEOREM 3.1. Let $\left\{B_{t}^{\delta}\right\}$ denote a Lévy-type process which has transition probability (3.2) with $x=0$ and $\mu=\sigma_{s}$. Then, up to a scale change, $\left\{B_{t}^{\delta}\right\}$ and $B E S^{\delta}(0)$ have the same distribution. Consequently, they are induced by the Kingman convolution.

Proof. Let $P_{x}^{\delta}$ denote the law of $B E S^{\delta}(x), \delta \geqslant 0, x \geqslant 0$, on $C\left(\mathbb{R}^{+}, \mathbb{R}\right)$ (cf. [12], XI, p. 446) which entails that the density $p_{t}^{\delta}(0, y)$ of the Bessel semigroup is

$$
\begin{equation*}
p_{t}^{\delta}(0, y)=2^{-s} t^{-(s+1)} \Gamma(s+1)^{-1} y^{2 s+1} \exp \left(-y^{2} / 2 t\right) \tag{3.3}
\end{equation*}
$$

It should be noted that functions (3.3) are Rayleigh functions of $y$. In addition, if $t=2$, we get $P_{2}^{\delta}(0, \cdot)=\sigma_{s}$. Next, by (1.10), we have

$$
\begin{equation*}
\widehat{\sigma_{s}^{\circ t}}(u)=\exp \left(-t u^{2} / 4(s+1)\right), \quad u \geqslant 0 \tag{3.4}
\end{equation*}
$$

Our further aim is to prove that, up to a scale change, the rad.ch.f. of $\sigma_{s}^{\circ t}$ is equal to the rad.ch.f. of $P_{t}^{\delta}(0, y)$. Accordingly, integrating the kernel $\Lambda_{s}(u z)$ with respect to $P_{t}^{\delta}(0, z)$ we see, by (1.4), (1.6) and (3.3), that the rad.ch.f. of $P_{t}^{\delta}(0, y)$ is given, for each $u \geqslant 0$, by

$$
\begin{align*}
\widehat{P_{t}^{\delta}}(0, y)(u) & =\int_{0}^{\infty} \Lambda_{s}(u z) P_{t}^{\delta}(0, z) d z  \tag{3.5}\\
& =2^{-s} t^{-(s+1)} \Gamma(s+1)^{-1} \int_{0}^{\infty} z^{2 s+1} \Lambda_{s}(u z) \exp \left(-z^{2} / 2 t\right) d z
\end{align*}
$$

Hence and by virtue of the Weber integral ${ }^{3}$ for $u \geqslant 0$ we have
(3.6) $\widehat{q_{t}^{\delta}}(0, y)(u)$

$$
\begin{aligned}
& =\left\{2^{-s} t^{-(s+1)} \Gamma(s+1)^{-1}\right\}\left\{2^{-1} 2^{s+1} t^{s+1} \Gamma(s+1) \exp \left(-t u^{2} / 2\right)\right\} \\
& =\widehat{\sigma_{s}^{\circ t}}(u)
\end{aligned}
$$

which shows that $q_{t}^{\delta}(0)=\sigma_{s}^{\circ t}$.
${ }^{3}$ From Watson ([18], p. 394) we have, for $s \geqslant-1 / 2, a \geqslant 0, p>0$,

$$
\int_{0}^{\infty} t^{s+1} J_{s}(a t) \exp \left(-p^{2} t^{2}\right) d t=a^{s}\left(2 p^{2}\right)^{-s-1} \exp \left(-a^{2} / 4 p^{2}\right)
$$

which may be written as

$$
\int_{0}^{\infty} t^{2 s+1} \Lambda_{s}(a t) \exp \left(-p^{2} t^{2}\right) d t=\frac{1}{2} \Gamma(s+1) p^{-2(s+1)} \exp \left(-a^{2} /\left(4 p^{2}\right)\right)
$$

## 4. BESSEL PROCESSES AS STATIONARY INDEPENDENT "INCREMENTS" PROCESSES

Suppose that $X_{j}, j=1,2, \ldots$, are nonnegative independent r.v.'s with the corresponding distributions $F_{X_{j}}, j=1,2, \ldots$, and $\theta, \theta_{1}, \theta_{2}, \ldots$ are i.i.d. r.v.'s with the common distribution $F_{s}$ and the r.v.'s $X_{j}, j=1,2, \ldots, \theta, \theta_{1}, \theta_{2}, \ldots$ are independent. Following Kingman ([5], formula (10)) we say, for a fixed $s \geqslant-1 / 2$, that any one of the equivalent r.v.'s

$$
\begin{equation*}
X_{1} \oplus X_{2}:=\sqrt{X_{1}^{2}+X_{2}^{2}+2 X_{1} X_{2} \theta_{1}} \tag{4.1}
\end{equation*}
$$

is a radial sum of the two independent nonnegative r.v.'s $X_{1}, X_{2}$. By induction, the radial sum $X_{1} \oplus X_{2} \oplus \ldots \oplus X_{k}$ is defined for any finite $k=2,3, \ldots$ It should be noted ([5], formula (12)) that the operation $\oplus$ is associative.

Definition 4.1. Let $\mathcal{B}_{b}$ be the ring of subsets of a non-empty bounded Borel subsets of $\mathbb{R}^{+}$. A function

$$
\begin{equation*}
M: \mathcal{B}_{b} \rightarrow L^{+} \tag{4.2}
\end{equation*}
$$

where $L^{+}=L^{+}(\Omega, \mathcal{F}, P)$ denotes the class of all nonnegative r.v.'s on the probability space $(\Omega, \mathcal{F}, P)$, is said to be an o-scattered random measure if
(i) $M(\emptyset)=0$ with probability 1 ;
(ii) for any $A, B \in \mathcal{B}_{b}, A \cap B=\emptyset$; then $M(A)$ and $M(B)$ are independent and

$$
\begin{equation*}
M(A \cup B) \stackrel{d}{=} M(A) \oplus M(B) \tag{4.3}
\end{equation*}
$$

(iii) for any pairwise disjoint sets $A_{1}, A_{2}, \ldots \in \mathcal{B}_{b}$ with the union in $\mathcal{B}_{b}$, the r.v.'s $M\left(A_{1}\right), M\left(A_{2}\right), \ldots$ are independent and

$$
\begin{equation*}
M\left(\bigcup_{j=1}^{\infty} A_{j}\right) \stackrel{d}{=} \bigoplus_{j=1}^{\infty} M\left(A_{j}\right) . \tag{4.4}
\end{equation*}
$$

It is well known that if $\{W(t)\}, t \in \mathbb{R}^{+}$, is a Brownian motion process, then there exists a Gaussian stochastic measure $M(A), A \in \mathcal{B}_{0}$, where $\mathcal{B}_{0}$ is the ring of bounded Borel subsets of $\mathbb{R}^{+}$with the property that, for every $t \geqslant 0$, we have $W(t)=M((0, t])$. The same is also true for Bessel processes. Namely, we get

ThEOREM 4.1. Let $\left\{B_{t}^{\delta}\right\}$ denote a Bessel process started at 0 . Then there exists a unique (up to finite-dimensional distributions) o-scattered random measure $B(A), A \in \mathcal{B}_{b}$, with the Lebesgue measure as its control measure such that for each $t \geqslant s \geqslant 0$ we have

$$
\begin{equation*}
B([0, t])=B_{s}^{\delta} \oplus B((s, t]) \tag{4.5}
\end{equation*}
$$

Moreover, the control measure associated with $B$ is the Lebesgue measure.

We proceed the proof of the theorem by showing the following lemma.
LEMMA 4.1. Let $\pi:=\left\{0=t_{0}<t_{1}<t_{2}<\ldots\right\}$ be a subdivision of $\mathbb{R}^{+}$. Then there exist independent r.v.'s $X_{1}, X_{2}, \ldots$ such that

$$
\begin{equation*}
X_{k} \stackrel{d}{=} \sigma_{s}^{t_{k}-t_{k-1}}, \quad k=0,1,2, \ldots \tag{4.6}
\end{equation*}
$$

Moreover, we have

$$
\begin{equation*}
B_{t_{n}}^{\delta} \stackrel{d}{=} X_{1} \oplus X_{2} \oplus \ldots \oplus X_{n}, \quad n=2,3, \ldots \tag{4.7}
\end{equation*}
$$

and

$$
\begin{equation*}
B\left(\left(t_{n}, t_{(n+r]}\right)\right) \stackrel{d}{=} \sigma_{s}^{t_{n+r}-t_{n}} \tag{4.8}
\end{equation*}
$$

Proof. Following the idea of Kingman ([5], p. 20) let us take as a sample space $\Omega$ the Cartesian product of countably many intervals $\mathbb{R}^{+}$with countably many intervals $[-1,1]$. The probability measure is defined on $\Omega$ as the product of the distributions $\sigma_{s}^{t_{k}-t_{k-1}}, k=1,2, \ldots$, on each of the first set of $\mathbb{R}^{+}$together with the distribution $F_{s}$ (cf. (1.8)) on each of the second set. If the typical point $\omega \in \Omega$ has components

$$
\begin{equation*}
X_{1}(\omega), X_{2}(\omega), \ldots ; \eta_{1}(\omega), \eta_{2}(\omega), \ldots \tag{4.9}
\end{equation*}
$$

define $S_{m}(\omega)$ inductively by

$$
\begin{gather*}
S_{0}=0  \tag{4.10}\\
S_{m+1}(\omega)=\left\{S_{m}^{2}(\omega)+X_{m+1}^{2}(\omega)+2 \eta_{m}(\omega) S_{m}(\omega) X_{m+1}(\omega)\right\}^{1 / 2} \tag{4.11}
\end{gather*}
$$

Thus, we have

$$
\begin{equation*}
S_{m+1}=S_{m} \oplus X_{m+1} \tag{4.12}
\end{equation*}
$$

which, by virtue of the associativity of $\oplus$, implies that for each $m=2,3, \ldots$

$$
\begin{equation*}
S_{m}=X_{1} \oplus X_{2} \oplus \ldots \oplus X_{m} \tag{4.13}
\end{equation*}
$$

Moreover, since $X_{k}, k=2,3, \ldots$, are independent, it follows that

$$
\begin{equation*}
S_{m} \stackrel{d}{=} \sigma^{t_{m}} \stackrel{d}{=} B\left(t_{m}\right) \tag{4.14}
\end{equation*}
$$

Now, since the operation $\oplus$ is associative (cf. Kingman [5], Theorem 1), one can show that

$$
\begin{gather*}
S_{m+r}=S_{m} \oplus S_{r}^{m} \\
S_{0}^{m}=0, \quad S_{r+1}^{m} X=S_{r}^{m} \oplus X_{m+r+1} \tag{4.15}
\end{gather*}
$$

Note, by (4.14) and (4.15), $\sigma^{t_{m+r}-t_{m}} \stackrel{d}{=} S_{r}^{m} \stackrel{d}{=}\left(X_{m} \oplus \ldots \oplus X_{m+r}\right)$, which entails (4.6), (4.7) and (4.8).

Proof of Theorem 4.1. Let $\mathbb{B}_{0}$ denote the class of finite unions of disjoint finite intervals ( $a, b$ ], i.e.

$$
\begin{equation*}
\bigcup_{j=1}^{k} I_{j}, I_{j}=\left(t_{2 j}, t_{2 j+1}\right], \quad j=0,1, \ldots, k=1,2, \ldots \tag{4.16}
\end{equation*}
$$

We put

$$
\begin{equation*}
B\left(\bigcup_{j=1}^{k} I_{j}\right)=\bigoplus_{j=1}^{k} B\left(I_{j}\right) . \tag{4.17}
\end{equation*}
$$

Using the transfinite induction, Lemma 4.1, and the usual extension method of random interval functions one can get an o-random measure $B(\cdot)$ on $\mathcal{B}_{b}$ with the required properties.

Definition 4.2. For every $0 \leqslant a \leqslant b$ the quantity $M((a, b])$ is called the increment-type of the Bessel processes $B E S^{\delta}(0)$.

Moreover, from (3.1) and (4.1) we have
Theorem 4.2. Every Bessel process which starts at 0 has a modification as a process with stationary and increments-type process.

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[^1]:    ${ }^{1}$ If $n$ is replaced by a negative real number, then the corresponding unique strong solution to the equation (1.3) exists, and thus the Bessel process of a negative dimension $\delta$ can be defined (cf. Revuz and Yor [12], Exercise 1.33, p. 453).

[^2]:    ${ }^{2}$ Higher dimensional Urbanik convolution algebras can be introduced in the same way as here for the Kingman convolution case but this subject will be treated systematically elsewhere.

